# IV.3. Stationary Markov Processes 

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Recap. Last lecture, we talked about two types of Markov processes: the Poisson process and the Brownian motion process. Both of these processes are lacking another property that can be useful in analyzing stochastic processes, that of stationarity, that we defined some time ago.

Stationarity and some notation. Recall from III.1: A stochastic process $Y$ is stationary if the moments are not affected by a time shift, i.e.,

$$
\left\langle Y\left(t_{1}+\tau\right) Y\left(t_{2}+\tau\right) \ldots Y\left(t_{n}+\tau\right)\right\rangle=\left\langle Y\left(t_{1}\right) Y\left(t_{2}\right) \ldots Y\left(t_{n}\right)\right\rangle
$$

for all $n, \tau$, and $t_{1}, t_{2}, \ldots, t_{n}$.
A theorem that applies only for Markov processes: A Markov process is stationary if and only if i) $P_{1}(y, t)$ does not depend on $t$; and ii) $P_{1 \mid 1}\left(y_{2}, t_{2} \mid y_{1}, t_{1}\right)$ depends only on the difference $t_{2}-t_{1}$. Condition ii) implies that $P_{1 \mid 1}\left(y_{2}, t_{2} \mid y_{1}, t_{1}\right)=P_{1 \mid 1}\left(y_{2}, t_{2}+\tau \mid y_{1}, t_{1}+\tau\right)$.

Proof: First suppose that conditions i) and ii) are satisfied. Then

$$
\begin{aligned}
\left\langle Y\left(t_{1}\right) Y\right. & \left.Y\left(t_{2}\right) \ldots Y\left(t_{n}\right)\right\rangle=\int y_{1} \ldots y_{n} P_{n}\left(y_{1}, t_{1} ; \ldots ; y_{n}, t_{n}\right) d y_{1} \ldots d y_{n} \\
& =\int y_{1} \ldots y_{n} P_{1 \mid 1}\left(y_{n}, t_{n} \mid y_{n-1}, t_{n-1}\right) \ldots P_{1 \mid 1}\left(y_{2}, t_{2} \mid y_{1}, t_{1}\right) P_{1}\left(y_{1}, t_{1}\right) d y_{1} \ldots d y_{n} \\
& =\int y_{1} \ldots y_{n} P_{1 \mid 1}\left(y_{n}, t_{n}+\tau \mid y_{n-1}, t_{n-1}+\tau\right) \ldots P_{1 \mid 1}\left(y_{2}, t_{2}+\tau \mid y_{1}, t_{1}+\tau\right) P_{1}\left(y_{1}, t_{1}+\tau\right) d y_{1} \ldots d y_{n} \\
& =\left\langle Y\left(t_{1}+\tau\right) Y\left(t_{2}+\tau\right) \ldots Y\left(t_{n}+\tau\right)\right\rangle
\end{aligned}
$$

Thus all moments are invariant under a time shift $\tau$.
Now suppose that the stationarity condition is satisfied. Specifically, this implies that $\left\langle Y^{n}(t)\right\rangle=\left\langle Y^{n}(t+\tau)\right\rangle$ for all $n$ and $\tau$. Since all the moments of $Y(t)$ and $Y(t+\tau)$ are equal, they must have the same probability distribution. Thus $P_{1}(y, t)=P_{1}(y, t+\tau)$ for all $\tau$, and thus it must not depend on $\tau$.

Consider the second moments of the process. If they are invariant under a time shift, it follows that

$$
\begin{aligned}
\left\langle Y\left(t_{1}\right) Y\left(t_{2}\right)\right\rangle & =\left\langle Y\left(t_{1}+\tau\right) Y\left(t_{2}+\tau\right)\right\rangle \\
\iint y_{1} y_{2} P_{2}\left(y_{1}, t_{1}, y_{2}, t_{2}\right) d y_{1} d y_{2} & =\iint y_{1} y_{2} P_{2}\left(y_{1}, t_{1}+\tau, y_{2}, t_{2}+\tau\right) d y_{1} d y_{2} \\
\iint y_{1} y_{2} P_{1 \mid 1}\left(y_{2}, t_{2} \mid y_{1}, t_{1}\right) P_{1}\left(y_{1}, t\right) d y_{1} d y_{2} & =\iint y_{1} y_{2} P_{1 \mid 1}\left(y_{2}, t_{2}+\tau \mid y_{1}, t_{1}+\tau\right) P_{1}\left(y_{1}, t\right) d y_{1} d y_{2}
\end{aligned}
$$

These two integrals can only be equal if $P_{1 \mid 1}\left(y_{2}, t_{2} \mid y_{1}, t_{1}\right)=P_{1 \mid 1}\left(y_{2}, t_{2}+\tau \mid y_{1}, t_{1}+\tau\right)$, which can only be guaranteed for all $\tau$ if $P_{1 \mid 1}\left(y_{2}, t_{2} \mid y_{1}, t_{1}\right)$ is a function of $t_{2}-t_{1}$. This completes the proof.

Therefore, we can define a more compact notation for stationary Markov processes:

$$
T_{\tau}\left(y_{2} \mid y_{1}\right):=P_{1 \mid 1}\left(y_{2}, t_{2} \mid y_{1}, t_{1}\right)
$$

This allows us to rewrite the Chapman-Kolmogorov Equation

$$
P_{1 \mid 1}\left(y_{3}, t_{3} \mid y_{1}, t_{1}\right)=\int P_{1 \mid 1}\left(y_{3}, t_{3} \mid y_{2}, t_{2}\right) P_{1 \mid 1}\left(y_{2}, t_{2} \mid y_{1}, t_{1}\right) d y_{2}
$$



Figure 1: Sample trajectory from a random telegraph process.
as

$$
T_{\tau+\tau^{\prime}}\left(y_{3} \mid y_{1}\right)=\int T_{\tau}\left(y_{3} \mid y_{2}\right) T_{\tau^{\prime}}\left(y_{2} \mid y_{1}\right) d y_{2}
$$

where $\tau^{\prime}:=t_{3}-t_{2}$. The Chapman-Kolmogorov Equation only applies when $\tau, \tau^{\prime}>0$.
Example. Suppose $y_{i}$ takes only integer values in the set $\{1,2,3, \ldots, n\}$. Then we can replace the integral with a sum to get

$$
T_{\tau+\tau^{\prime}}\left(y_{3} \mid y_{1}\right)=\sum_{y_{2}=1}^{n} T_{\tau}\left(y_{3} \mid y_{2}\right) T_{\tau^{\prime}}\left(y_{2} \mid y_{1}\right) .
$$

This is just the formula for finding each element of a matrix multiplication. So you can think of the ChapmanKolmogorov equation as being a matrix identity:

$$
\mathbf{T}_{\tau+\tau^{\prime}}=\mathbf{T}_{\tau} \mathbf{T}_{\tau^{\prime}} .
$$

If the state space is not finite, we can extend this idea from matrices to "integral kernels" in a similar fashion, resulting in the same equation.

Example: Random Telegraph Process. The random telegraph process is defined as a Markov process that takes on only two values: 1 and -1 , which it switches between with the rate $\gamma$. It can be defined by the equation

$$
\frac{\partial}{\partial t} P_{1}(y, t)=-\gamma P_{1}(y, t)+\gamma P_{1}(-y, t)
$$

When the process starts at $t=0$, it is equally likely that the process takes either value, that is

$$
P_{1}(y, 0)=\frac{1}{2} \delta(y-1)+\frac{1}{2} \delta(y+1) .
$$

Goal: To show that the random telegraph process is stationary. We'll need to show that $P_{1}(y, t)$ does not depend on $t$ and that $P_{1 \mid 1}\left(y_{2}, y_{1} \mid t_{2}, t_{1}\right)$ is a function of $t_{2}-t_{1}$.

How is the number of times that a given trajectory of the process switches between 1 and -1 in a given interval $\left(t_{1}, t_{2}\right]$ ? Recall from Monday's lecture, the Poisson process. The process is like a Poisson process except that instead of increasing by 1 each time a new arrival occurs, it switches. We can prove by induction that the distribution of arrivals in any intervals $\left(t_{1}, t_{2}\right]$ is Poisson.

Base step: The probability there are no switches in an interval $\left(t^{\prime}, t^{\prime}+d t^{\prime}\right]$ is $1-\gamma d t^{\prime}$ for small $d t^{\prime}$. The probability that there are no switches in $\left(t_{1}, t_{2}\right]$ is then

$$
\operatorname{Pr}\left(0 \text { switches in }\left(t_{1}, t_{2}\right]\right)=\lim _{d t^{\prime} \rightarrow 0}\left(1-\gamma d t^{\prime}\right)^{\frac{t_{2}-t_{1}}{d t^{\prime}}}=e^{-\gamma\left(t_{2}-t_{1}\right)}=e^{-\gamma\left(t_{2}-t_{1}\right)} \frac{\left(-\gamma\left(t_{2}-t_{1}\right)\right)^{0}}{0!}
$$

Induction Step: Assume that the probability of $n$ switches in the interval $\left(t_{1}, t_{2}\right]$ is $p_{n}=e^{-\gamma\left(t_{2}-t_{1}\right)} \frac{\left(\gamma\left(t_{2}-t_{1}\right)\right)^{n}}{n!}$ for $n=0 \ldots N$. Then to find the probability that there are $N+1$ switches in the interval, condition on the time of the 1st switch in the interval, which occurs at time $t^{\prime}$ with probability $\gamma d t^{\prime}$. Then there must be 0 switches in the interval $\left(t_{1}, t^{\prime}\right]$ and $N$ switches in the interval $\left(t^{\prime}, t_{2}\right]$. The probability of this is

$$
\begin{aligned}
\operatorname{Pr}\left(\mathrm{N}+1 \text { switches in }\left(t_{1}, t_{2}\right]\right) & =\int_{t_{1}}^{t_{2}} e^{-\gamma\left(t_{2}-t^{\prime}\right)} \frac{\left(\gamma\left(t_{2}-t^{\prime}\right)\right)^{n}}{n!} e^{-\gamma\left(t^{\prime}-t_{1}\right)} \gamma d t^{\prime} \\
& =\frac{e^{-\gamma\left(t_{2}-t_{1}\right)}}{n!} \gamma^{n+1} \int_{t_{1}}^{t_{2}}\left(t_{2}-t^{\prime}\right)^{n} d t^{\prime} \\
& =\frac{e^{-\gamma\left(t_{2}-t_{1}\right)}}{n!} \gamma^{n+1} \frac{\left(t_{2}-t_{1}\right)^{n+1}}{n+1} \\
& =\frac{e^{-\gamma\left(t_{2}-t_{1}\right)}}{(n+1)!}\left[\gamma\left(t_{2}-t_{1}\right)\right]^{n+1}
\end{aligned}
$$

Now let's find $P_{1 \mid 1}\left(y_{2}, y_{1} \mid t_{2}, t_{1}\right)$. If the trajectory generated by the process switches an even number of times, then $y_{1}=y_{2}$. If it switches an odd number of times, then $y_{1}=-y_{2}$. Therefore

$$
P_{1 \mid 1}\left(y_{2}, t_{2} \mid y_{1}, t_{1}\right)=\sum_{n=0,2,4, \ldots} e^{-\gamma\left(t_{2}-t_{1}\right)} \frac{\left(\gamma\left(t_{2}-t_{1}\right)\right)^{n}}{n!} \delta\left(y_{1}-y_{2}\right)+\sum_{n=1,3,5, \ldots} e^{-\gamma\left(t_{2}-t_{1}\right)} \frac{\left(\gamma\left(t_{2}-t_{1}\right)\right)^{n}}{n!} \delta\left(y_{1}+y_{2}\right)
$$

The power series in the two terms are those of hyperbolic cosine and hyperbolic sine, respectively, so

$$
\begin{aligned}
P_{1 \mid 1}\left(y_{2}, t_{2} \mid y_{1}, t_{1}\right) & =e^{-\gamma\left(t_{2}-t_{1}\right)} \cosh \left(\gamma\left(t_{2}-t_{1}\right)\right) \delta\left(y_{1}-y_{2}\right)+e^{-\gamma\left(t_{2}-t_{1}\right)} \sinh \left(\gamma\left(t_{2}-t_{1}\right)\right) \delta\left(y_{1}+y_{2}\right) \\
& =e^{-\gamma\left(t_{2}-t_{1}\right)}\left(\frac{e^{\gamma\left(t_{2}-t_{1}\right)}+e^{-\gamma\left(t_{2}-t_{1}\right)}}{2}\right) \delta\left(y_{1}-y_{2}\right)+e^{-\gamma\left(t_{2}-t_{1}\right)}\left(\frac{e^{\gamma\left(t_{2}-t_{1}\right)}-e^{-\gamma\left(t_{2}-t_{1}\right)}}{2}\right) \delta\left(y_{1}+y_{2}\right) \\
& =\frac{1}{2}\left(1+e^{-2 \gamma\left(t_{2}-t_{1}\right)}\right) \delta\left(y_{2}-y_{1}\right)+\frac{1}{2}\left(1-e^{-2 \gamma\left(t_{2}-t_{1}\right)}\right) \delta\left(y_{2}+y_{1}\right)
\end{aligned}
$$

The probability distribution for any time $t$ can be found as follows:

$$
\begin{aligned}
P_{1}\left(y_{2}, t\right) & =P_{1 \mid 1}\left(y_{2}, t \mid y_{1}, 0\right) P_{1}\left(y_{1}, 0\right) \\
& =\left(\frac{1}{2}\left(1+e^{-2 \gamma t_{2}}\right) \delta\left(y_{2}-y_{1}\right)+\frac{1}{2}\left(1-e^{-2 \gamma t_{2}}\right) \delta\left(y_{2}+y_{1}\right)\right)\left(\frac{1}{2} \delta\left(y_{1}-1\right)+\frac{1}{2} \delta\left(y_{1}+1\right)\right) \\
& =\frac{1}{4}\left(1+e^{-2 \gamma t_{2}}\right) \delta\left(y_{2}-1\right)+\frac{1}{4}\left(1-e^{-2 \gamma t_{2}}\right) \delta\left(y_{2}-1\right)+\frac{1}{4}\left(1+e^{-2 \gamma t_{2}}\right) \delta\left(y_{2}+1\right)+\frac{1}{4}\left(1-e^{-2 \gamma t_{2}}\right) \delta\left(y_{2}+1\right) \\
& =\frac{1}{2} \delta\left(y_{2}-1\right)+\frac{1}{2} \delta\left(y_{2}+1\right)
\end{aligned}
$$

Is the random telegraph process stationary? Yes, because, i) $P_{1}(y, t)$ does not depend on $t$, and ii) $P_{1 \mid 1}\left(y_{2}, y_{1} \mid t_{2}, t_{1}\right)$ is a function of $\tau=t_{2}-t_{1}$. Therefore we can write

$$
T_{\tau}\left(y_{2} \mid y_{1}\right)=\frac{1}{2}\left(1+e^{-2 \gamma \tau}\right) \delta\left(y_{2}-y_{1}\right)+\frac{1}{2}\left(1-e^{-2 \gamma \tau}\right) \delta\left(y_{2}+y_{1}\right)
$$

Autocorrelation of a stationary process. Since a stationary process has the same probability distribution for all time $t$, we can always shift the values of the $y$ 's by a constant to make the process a zero-mean process. So let's just assume $\langle Y(t)\rangle=0$. The autocorrelation function is thus:

$$
\kappa\left(t_{1}, t_{1}+\tau\right)=\left\langle Y\left(t_{1}\right) Y\left(t_{1}+\tau\right)\right\rangle
$$

Since the process is stationary, this doesn't depend on $t_{1}$, so we'll denote it by $\kappa(\tau)$. If we know expressions of the transition probability function and the unconditional probability function, we can calculate the autocorrelation function


Figure 2: Sample trajectory of the Ornstein-Uhlenbeck process. The dashed line is the integral of the trajectory, which should behave similarly to Brownian motion.
using the formula derived as follows.

$$
\begin{aligned}
\kappa(\tau) & =\iint y_{1} y_{2} P_{2}\left(y_{1}, t_{1}, y_{2}, t_{1}+\tau\right) d y_{1} d y_{2} \\
& =\iint y_{1} y_{2} P_{1 \mid 1}\left(y_{2}, t_{1}+\tau \mid y_{1}, t_{1}\right) P_{1}\left(y_{1}\right) d y_{1} d y_{2} \\
& =\iint y_{1} y_{2} T_{\tau}\left(y_{2} \mid y_{1}\right) P_{1}\left(y_{1}\right) d y_{1} d y_{2} .
\end{aligned}
$$

Example. Autocorrelation of the random telegraph process.

$$
\begin{aligned}
\kappa(\tau) & =\sum_{y_{1} \in\{-1,1\}} \sum_{y_{2} \in\{-1,1\}} y_{1} y_{2} T_{\tau}\left(y_{2} \mid y_{1}\right) P_{1}\left(y_{1}\right) \\
& =(1)(1) T_{\tau}(1 \mid 1) P_{1}(1)+(-1)(1) T_{\tau}(-1 \mid 1) P_{1}(1)+(1)(-1) T_{\tau}(1 \mid-1) P_{1}(-1)+(-1)(-1) T_{\tau}(-1 \mid-1) P_{1}(-1) \\
& =\frac{1}{2} \frac{1}{2}\left(1+e^{-2 \gamma \tau}\right)-\frac{1}{2} \frac{1}{2}\left(1-e^{-2 \gamma \tau}\right)-\frac{1}{2} \frac{1}{2}\left(1-e^{-2 \gamma \tau}\right)+\frac{1}{2} \frac{1}{2}\left(1+e^{-2 \gamma \tau}\right) \\
& =\frac{1}{2}\left(1+e^{-2 \gamma \tau}\right)-\frac{1}{2}\left(1-e^{-2 \gamma \tau}\right) \\
& =e^{-2 \gamma \tau}
\end{aligned}
$$

The Ornstein-Uhlenbeck Process. The Ornstein-Uhlenbeck process was constructed in order to describe the velocity of a particle in the physical process of Brownian motion. The Ornstein-Uhlenbeck process is a mathematically distinct entity for the Wiener-Levy process that describes the position of a particle in Brownian motion; you can't just integrate and differentiate between the two. It is a stationary Markov process defined by the following equations.

$$
\begin{aligned}
P_{1}\left(y_{1}\right) & =\frac{1}{\sqrt{2 \pi}} e^{-\frac{1}{2} y_{1}^{2}} \sim \mathcal{N}(0,1) \\
T_{\tau}\left(y_{2} \mid y_{1}\right) & =\frac{1}{\sqrt{2 \pi\left(1-e^{-2 \tau}\right)}} e^{-\frac{\left(y_{2}-y_{1} e^{-\tau}\right)^{2}}{2\left(1-e^{-2 \tau}\right)}} \sim \mathcal{N}\left(y_{1} e^{-\tau}, 1-e^{-2 \tau}\right) .
\end{aligned}
$$

For this process to be properly defined, the functions $P_{1}$ and $T_{\tau}$ must satisfy 1) the Chapman-Kolmogorov equation and 2) the consistency condition $\int T_{\tau}\left(y_{2} \mid y_{1}\right) P_{1}\left(y_{1}\right) d y_{1}=P_{1}\left(y_{2}\right)$. The Ornstein-Uhlenbeck process satisfies condition
2) as shown below:

$$
\begin{aligned}
\int T_{\tau}\left(y_{2} \mid y_{1}\right) P_{1}\left(y_{1}\right) d y_{1} & =\frac{1}{\sqrt{2 \pi}} \int \frac{1}{\sqrt{2 \pi\left(1-e^{-2 \tau}\right)}} \exp \left[-\frac{1}{2} y_{1}^{2}-\frac{\left(y_{2}-y_{1} e^{-\tau}\right)^{2}}{2\left(1-e^{-2 \tau}\right)}\right] d y_{1} \\
& =\frac{1}{\sqrt{2 \pi}} \int \frac{1}{\sqrt{2 \pi\left(1-e^{-2 \tau}\right)}} \exp \left[-\frac{\left(1-e^{-2 \tau}\right) y_{1}^{2}+y_{2}^{2}-2 y_{1} y_{2} e^{-\tau}+y_{1}^{2} e^{-2 \tau}}{2\left(1-e^{-2 \tau}\right)}\right] d y_{1} \\
& =\frac{1}{\sqrt{2 \pi}} \underbrace{\int \frac{1}{\sqrt{2 \pi\left(1-e^{-2 \tau}\right)}} \exp \left[-\frac{y_{1}^{2}-2 y_{1} y_{2} e^{-\tau}+y_{2}^{2} e^{-2 \tau}}{2\left(1-e^{-2 \tau}\right)}\right] d y_{1}}_{=1 \text { (Gaussian pdf) }} \exp \left[-\frac{1}{2} y_{2}^{2}\right] \\
& =\frac{1}{\sqrt{2 \pi}} e^{-\frac{1}{2} y_{2}^{2}}
\end{aligned}
$$

The Ornstein-Uhlenbeck process also satisfies the Chapman-Kolmogorov equation. The book states that "the reader will have no difficulty in verifying" that these conditions are satisfied. Conceptually, it's not difficult, but it is extremely tedious and skippable. Instead, let's find the autocorrelation function of the process:

$$
\begin{aligned}
\kappa(\tau) & =\int_{y_{1}} \int_{y_{2}} y_{1} P_{1}\left(y_{1}\right) y_{2} T_{\tau}\left(y_{2} \mid y_{1}\right) d y_{1} d y_{2} \\
& =\int_{y_{1}} y_{1} \frac{1}{\sqrt{2 \pi}} e^{-\frac{1}{2} y_{1}^{2}} \underbrace{\int_{y_{2}} y_{2} \frac{1}{\sqrt{2 \pi\left(1-e^{-2 \tau}\right)}} e^{-\frac{\left(y_{2}-y_{1} e^{-\tau}\right)^{2}}{2\left(1-e^{-2 \tau}\right)}} d y_{2}}_{=y_{1} e^{-\tau} \text { (mean of Gaussian pdf) }} d y_{1} \\
& =e^{-\tau} \underbrace{\int_{y_{1}} y_{1}^{2} \frac{1}{\sqrt{2 \pi}} e^{-\frac{1}{2} y_{1}^{2}} d y_{1}}_{=1 \text { (variance of Gaussian pdf) }} \\
& =e^{-\tau}
\end{aligned}
$$

Notice the the autocorrelation function of the Ornstein-Uhlenbeck process is the same form as that of the random telegraph process.

The Ornstein-Uhlenbeck process is interesting because it is essentially the only process that is Gaussian, Markov, and stationary. (Essentially means that processes that are translated in time or space are considered to be the same process, and one pathological process is excluded.) This result is called Doob's Theorem. The random telegraph process has only two of these properties: it's Markovian and stationary, but not Gaussian.

